

Intro to Quantum Computing  
CSCI/PHYS 3090  
CU Boulder Spring 2020  
Problem Set 3 Solutions

### Problem 1

We want to calculate the probabilities of getting  $|0\rangle$  and  $|1\rangle$  if we measure qubit one on the following states. We additionally would like to determine the post-measurement state on qubit 2 for each of the measurement outcomes.

The measurement probabilities and the post-measurement states can be determined by using the generalized Born-rule. We would like to get the states into the following form:

$$|\psi\rangle = \alpha_0 |0\rangle_1 \otimes |\psi_0\rangle_2 + \alpha_1 |1\rangle_1 \otimes |\psi_1\rangle_2 \quad (1)$$

where  $|\psi_0\rangle_2, |\psi_1\rangle_2$  are normalized states. If our state is in the form shown above, then we get outcome “x” with probability  $p_x = |\alpha_x|^2$  and the post-measurement state on qubit 2 is  $|\psi_x\rangle_2$ .

#### Part A

$$|\psi\rangle_{12} = \frac{1}{\sqrt{3}} |0\rangle_1 |1\rangle_2 + \sqrt{\frac{2}{3}} |1\rangle_1 |0\rangle_2$$

This state is already in the form given by Equation 1 and so we can get the following probabilities and post-measurement outcomes:

| Outcome | Probability   | Post-measurement state on qubit 2 |
|---------|---------------|-----------------------------------|
| 0       | $\frac{1}{3}$ | $ 1\rangle_2$                     |
| 1       | $\frac{2}{3}$ | $ 0\rangle_2$                     |

Table 1: Outcomes for Part A

## Part B

$$\begin{aligned}
|\psi\rangle_{12} &= \frac{1}{\sqrt{2}} |+\rangle_1 |0\rangle_2 + \frac{1}{\sqrt{2}} |-\rangle_1 |1\rangle_2 \\
|\psi\rangle_{12} &= \frac{1}{2}(|0\rangle_1 + |1\rangle_1) |0\rangle_2 + \frac{1}{2}(|0\rangle_1 - |1\rangle_1) |1\rangle_2 \\
|\psi\rangle_{12} &= \frac{1}{2}(|0\rangle_1 |0\rangle_2 + |0\rangle_1 |1\rangle_2 + |1\rangle_1 |0\rangle_2 - |1\rangle_1 |1\rangle_2) \\
|\psi\rangle_{12} &= \frac{1}{2}(|0\rangle_1 (|0\rangle_2 + |1\rangle_2) + |1\rangle_1 (|0\rangle_2 - |1\rangle_2)) \\
|\psi\rangle_{12} &= \frac{1}{\sqrt{2}}(|0\rangle_1 |+\rangle_2 + |1\rangle_1 |-\rangle_2)
\end{aligned}$$

Therefore we get the following outcomes:

| Outcome | Probability   | Post-measurement state on qubit 2 |
|---------|---------------|-----------------------------------|
| 0       | $\frac{1}{2}$ | $ +\rangle_2$                     |
| 1       | $\frac{1}{2}$ | $ -\rangle_2$                     |

Table 2: Outcomes for Part B

## Part C

$$\begin{aligned}
|\psi\rangle_{12} &= \frac{1}{\sqrt{2+\sqrt{2}}}(|+\rangle_1 |0\rangle_2 + |+\rangle_1 |-\rangle_2) \\
|\psi\rangle_{12} &= \frac{1}{\sqrt{2+\sqrt{2}}} |+\rangle_1 (|0\rangle_2 + |-\rangle_2) \\
|\psi\rangle_{12} &= \frac{1}{\sqrt{2+\sqrt{2}}} \frac{1}{\sqrt{2}} (|0\rangle_1 + |1\rangle_1) (|0\rangle_2 + |-\rangle_2)
\end{aligned}$$

We now want to normalize the second qubit state in order to properly apply the generalized Born rule.

$$\begin{aligned}
|\psi\rangle_{12} &= \frac{1}{\sqrt{2+\sqrt{2}}} \frac{1}{\sqrt{2}} (|0\rangle_1 + |1\rangle_1) (|0\rangle_2 + \frac{1}{\sqrt{2}} (|0\rangle_2 - |1\rangle_2)) \\
|\psi\rangle_{12} &= \frac{1}{\sqrt{2+\sqrt{2}}} \frac{1}{\sqrt{2}} (|0\rangle_1 + |1\rangle_1) ((1 + \frac{1}{\sqrt{2}}) |0\rangle_2 - \frac{1}{\sqrt{2}} |1\rangle_2) \\
\sqrt{(1 + \frac{1}{\sqrt{2}})^2 + (\frac{1}{\sqrt{2}})^2} &= \sqrt{1 + \frac{2}{\sqrt{2}} + \frac{1}{2} + \frac{1}{2}} = \sqrt{2 + \sqrt{2}} \\
|\psi\rangle_{12} &= \frac{1}{\sqrt{2}} (|0\rangle_1 + |1\rangle_1) (\frac{1 + \frac{1}{\sqrt{2}}}{\sqrt{2+\sqrt{2}}} |0\rangle_2 - \frac{1}{\sqrt{2}\sqrt{2+\sqrt{2}}} |1\rangle_2) \\
|\psi\rangle_{12} &= \frac{1}{\sqrt{2}} (|0\rangle_1 + |1\rangle_1) (\frac{1 + \sqrt{2}}{\sqrt{4+2\sqrt{2}}} |0\rangle_2 - \frac{1}{\sqrt{4+2\sqrt{2}}} |1\rangle_2)
\end{aligned}$$

We now have the state in a form where we can simply use the Born rule to get the measurement outcomes.

| Outcome | Probability   | Post-measurement state on qubit 2                                                              |
|---------|---------------|------------------------------------------------------------------------------------------------|
| 0       | $\frac{1}{2}$ | $\frac{1+\sqrt{2}}{\sqrt{4+2\sqrt{2}}}  0\rangle_2 - \frac{1}{\sqrt{4+2\sqrt{2}}}  1\rangle_2$ |
| 1       | $\frac{1}{2}$ | $\frac{1+\sqrt{2}}{\sqrt{4+2\sqrt{2}}}  0\rangle_2 - \frac{1}{\sqrt{4+2\sqrt{2}}}  1\rangle_2$ |

Table 3: Outcomes for Part C

## Part D

$$\begin{aligned}
|\psi\rangle_{12} &= \frac{1}{\sqrt{2}}(|0\rangle_1 |+\rangle_2 + |+\rangle_1 |-\rangle_2) \\
|\psi\rangle_{12} &= \frac{1}{\sqrt{2}}(|0\rangle_1 |+\rangle_2 + \frac{1}{\sqrt{2}}(|0\rangle_1 + |1\rangle_1) |-\rangle_2) \\
|\psi\rangle_{12} &= |0\rangle_1 \left(\frac{1}{\sqrt{2}} |+\rangle_2 + \frac{1}{2} |-\rangle_2\right) + \frac{1}{2} |1\rangle_1 |-\rangle_2 \\
|\psi\rangle_{12} &= \sqrt{\frac{3}{4}} |0\rangle_1 \left(\sqrt{\frac{2}{3}} |+\rangle_2 + \frac{1}{\sqrt{3}} |-\rangle_2\right) + \frac{1}{2} |1\rangle_1 |-\rangle_2
\end{aligned}$$

| Outcome | Probability   | Post-measurement state on qubit 2                                 |
|---------|---------------|-------------------------------------------------------------------|
| 0       | $\frac{3}{4}$ | $\sqrt{\frac{2}{3}}  +\rangle_2 + \frac{1}{\sqrt{3}}  -\rangle_2$ |
| 1       | $\frac{1}{4}$ | $ -\rangle_2$                                                     |

Table 4: Outcomes for Part D

## Problem 2

Given that  $|u\rangle, |v\rangle, |w\rangle$ , and  $|z\rangle$  are real vectors, then we want to show that the inner product of  $(|u \otimes v\rangle, |w \otimes z\rangle)$  is equal to  $\langle u|w\rangle \cdot \langle v|z\rangle$ . We will have to assume that  $|u\rangle$  and  $|w\rangle$  have the same dimension and that  $|v\rangle$  and  $|z\rangle$  have the same dimension for the inner product to be well defined.

We can define the tensor product in terms of how it operates on vectors in some basis:

$$u = \begin{pmatrix} u_1 \\ u_2 \\ \vdots \\ u_m \end{pmatrix}, \quad v = \begin{pmatrix} v_1 \\ v_2 \\ \vdots \\ v_n \end{pmatrix}, \quad w = \begin{pmatrix} w_1 \\ w_2 \\ \vdots \\ w_m \end{pmatrix}, \quad z = \begin{pmatrix} z_1 \\ z_2 \\ \vdots \\ z_n \end{pmatrix}$$

$$|u\rangle \otimes |v\rangle = \begin{pmatrix} u_1v_1 \\ u_1v_2 \\ \vdots \\ u_1v_n \\ u_2v_1 \\ \vdots \\ u_mv_n \end{pmatrix}, \quad |w\rangle \otimes |z\rangle = \begin{pmatrix} w_1z_1 \\ w_1z_2 \\ \vdots \\ w_1z_n \\ w_2z_1 \\ \vdots \\ w_mz_n \end{pmatrix}$$

The inner product is the following:

$$\begin{aligned} (|u\rangle \otimes |v\rangle, |w\rangle \otimes |z\rangle) &= u_1v_1w_1z_1 + u_1v_2w_1z_2 + \cdots + u_1v_nw_1z_n + u_2v_1w_2z_1 + \cdots + u_mv_nw_mz_n \\ (|u\rangle \otimes |v\rangle, |w\rangle \otimes |z\rangle) &= (u_1w_1)(v_1z_1 + v_2z_2 + \cdots + v_nz_n) + \cdots + (u_mv_m)(v_1z_1 + v_2z_2 + \cdots + v_nz_n) \\ (|u\rangle \otimes |v\rangle, |w\rangle \otimes |z\rangle) &= (u_1w_1 + u_2w_2 + \cdots + u_mw_m) \langle v|z\rangle \\ (|u\rangle \otimes |v\rangle, |w\rangle \otimes |z\rangle) &= \langle u|w\rangle \cdot \langle v|z\rangle \end{aligned}$$

### Problem 3

We begin by defining the bitwise inner product between two bit strings  $x = (x_1, \dots, x_n)$  and  $a = (a_1, \dots, a_n)$  as the following:

$$x \cdot a = x_1a_1 + \cdots + x_na_n$$

where addition is assumed to mean addition modulo 2. The addition between two bit strings in this problem is bitwise addition modulo 2:  $x + a = (a_1 + x_1, \dots, a_n + x_n)$ . We would like to prove the following relation:

$$\sum_{x \in \{0,1\}^n} (-1)^{(a+y) \cdot x} = \prod_{j=1}^n \sum_{x_j=0}^1 (-1)^{(a_j+y_j)x_j} \quad (2)$$

### Part A

We will begin by considering the  $n = 1$  case. Let's begin by looking at the RHS.

$$\prod_{j=1}^1 \sum_{x_j=0}^1 (-1)^{(a_j+y_j)x_j} = \sum_{x_1=0}^1 (-1)^{(a_1+y_1)x_1} = (-1)^{(a_1+y_1) \cdot 0} + (-1)^{(a_1+y_1) \cdot 1}$$

We can now consider the LHS:

$$\begin{aligned} \sum_{x \in \{0,1\}^1} (-1)^{(a+y) \cdot x} &= (-1)^{(a+y) \cdot 0} + (-1)^{(a+y) \cdot 1} = (-1)^{(a_1+y_1) \cdot 0} + (-1)^{(a_1+y_1) \cdot 1} \\ \sum_{x \in \{0,1\}^1} (-1)^{(a+y) \cdot x} &= (-1)^{(a_1+y_1) \cdot 0} + (-1)^{(a_1+y_1) \cdot 1} \end{aligned}$$

We therefore see that the Equation 2 holds for the case of  $n = 1$ .

## Part B

Now let's consider the case of  $n = 2$ . We will try to work the RHS of the equation into the LHS.

$$\begin{aligned}
\prod_{j=1}^2 \sum_{x_j=0}^1 (-1)^{(a_j+y_j)x_j} &= \prod_{j=1}^2 ((-1)^{(a_j+y_j)0} + (-1)^{(a_j+y_j)1}) \\
\prod_{j=1}^2 \sum_{x_j=0}^1 (-1)^{(a_j+y_j)x_j} &= ((-1)^{(a_1+y_1)0} + (-1)^{(a_1+y_1)1})((-1)^{(a_2+y_2)0} + (-1)^{(a_2+y_2)1}) \\
\prod_{j=1}^2 \sum_{x_j=0}^1 (-1)^{(a_j+y_j)x_j} &= (-1)^{(a_1+y_1)0}(-1)^{(a_2+y_2)0} + (-1)^{(a_1+y_1)1}(-1)^{(a_2+y_2)0} \\
&\quad + (-1)^{(a_1+y_1)0}(-1)^{(a_2+y_2)0} + (-1)^{(a_1+y_1)1}(-1)^{(a_2+y_2)1} \\
\prod_{j=1}^2 \sum_{x_j=0}^1 (-1)^{(a_j+y_j)x_j} &= (-1)^{(a_1+y_1)0+(a_2+y_2)0} + (-1)^{(a_1+y_1)1+(a_2+y_2)0} \\
&\quad + (-1)^{(a_1+y_1)0+(a_2+y_2)1} + (-1)^{(a_1+y_1)1+(a_2+y_2)1} \\
\prod_{j=1}^2 \sum_{x_j=0}^1 (-1)^{(a_j+y_j)x_j} &= (-1)^{(a+y)\cdot(0,0)} + (-1)^{(a+y)\cdot(1,0)} + (-1)^{(a+y)\cdot(0,1)} + (-1)^{(a+y)\cdot(1,1)} \\
\prod_{j=1}^2 \sum_{x_j=0}^1 (-1)^{(a_j+y_j)x_j} &= \sum_{x \in \{0,1\}^2} (-1)^{(a+y)\cdot x}
\end{aligned}$$

Therefore, we have shown that Equation 2 holds for the  $n = 2$  case.

## Part C

We will now show the inductive step by showing that if we assume that the identity holds for the  $n$ -th case, then it holds for the  $(n+1)$ -th case. We will start with the RHS and prove that it is equal to the LHS. Define the following strings:  $a' = (a_1, \dots, a_n)$  and  $y' = (y_1, \dots, y_n)$  where the bit strings  $a$  and  $y$  that we are given are  $(a', a_{n+1})$  and  $(y', y_{n+1})$ .

$$\begin{aligned}
\prod_{j=1}^{n+1} \sum_{x_j=0}^1 (-1)^{(a_j+y_j)x_j} &= \left( \sum_{x_{n+1}=0}^1 (-1)^{(a_{n+1}+y_{n+1})x_{n+1}} \right) \prod_{j=1}^n \sum_{x_j=0}^1 (-1)^{(a_j+y_j)x_j} \\
\prod_{j=1}^{n+1} \sum_{x_j=0}^1 (-1)^{(a_j+y_j)x_j} &= ((-1)^{(a_{n+1}+y_{n+1})0} + (-1)^{(a_{n+1}+y_{n+1})1}) \sum_{x \in \{0,1\}^n} (-1)^{(a'+y')\cdot x} \\
\prod_{j=1}^{n+1} \sum_{x_j=0}^1 (-1)^{(a_j+y_j)x_j} &= \sum_{x \in \{0,1\}^n} (-1)^{(a_{n+1}+y_{n+1})0} (-1)^{(a'+y')\cdot x} + \sum_{x \in \{0,1\}^n} (-1)^{(a_{n+1}+y_{n+1})1} (-1)^{(a'+y')\cdot x} \\
\prod_{j=1}^{n+1} \sum_{x_j=0}^1 (-1)^{(a_j+y_j)x_j} &= \sum_{x \in \{0,1\}^n} (-1)^{(a'+y')\cdot x + (a_{n+1}+y_{n+1})0} + \sum_{x \in \{0,1\}^n} (-1)^{(a'+y')\cdot x + (a_{n+1}+y_{n+1})1} \\
\prod_{j=1}^{n+1} \sum_{x_j=0}^1 (-1)^{(a_j+y_j)x_j} &= \sum_{x \in \{0,1\}^n} (-1)^{((a', a_{n+1}) + (y', y_{n+1})) \cdot (x, 0)} + \sum_{x \in \{0,1\}^n} (-1)^{((a', a_{n+1}) + (y', y_{n+1})) \cdot (x, 1)} \\
\prod_{j=1}^{n+1} \sum_{x_j=0}^1 (-1)^{(a_j+y_j)x_j} &= \sum_{x \in \{0,1\}^n} (-1)^{(a+y)\cdot(x, 0)} + \sum_{x \in \{0,1\}^n} (-1)^{(a+y)\cdot(x, 1)}
\end{aligned}$$

We now observe that we are really summing over all possible bit strings of length  $n + 1$  since we are adding together the case where the  $(n + 1)$ -th bit is 0 and the case when it is 1 for each possible string of  $n$  least significant bits.

$$\prod_{j=1}^{n+1} \sum_{x_j=0}^1 (-1)^{(a_j+y_j)x_j} = \sum_{x \in \{0,1\}^{n+1}} (-1)^{(a+y) \cdot x}$$

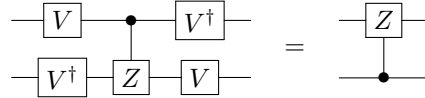
Therefore we have shown the inductive step and proven by induction that Equation 2 holds.

## Problem 4

In this problem we are trying to find unitaries that will interchange the control and target qubits for ControlZ and ControlY gates.

### Part A

We want to find a unitary  $V$  such that the following is true:



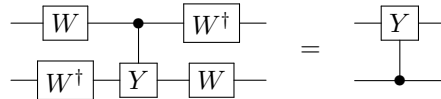
We can start by just expressing  $CZ_{12}$  and  $CZ_{21}$  as 4x4 matrices.

$$CZ_{12} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}$$

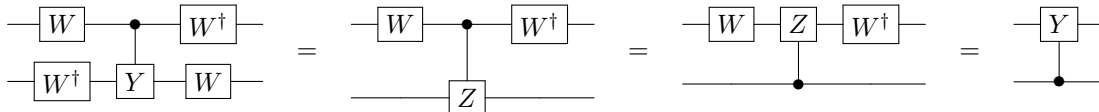
$$CZ_{21} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}$$

The  $CZ$  gate is symmetric under exchange of the control and target qubits as the only computational basis state that the gate acts on non-trivially is  $|1\rangle_1 |1\rangle_2$ . Therefore the unitary that satisfies the above circuit equation is  $V = I$  (the identity).

### Part B



We can determine the unitary  $W$  by noting that in the previous problem we found that the  $CZ$  gate is symmetric in its target and control qubits. This means that if we can find a gate  $W$  such that  $W^\dagger Y W = Z$ , then we can turn  $CY_{12}$  into  $CZ_{12}$ . This is equivalent to  $CZ_{21}$ , so then when we act the inverse gates that take us back from  $Z$  to  $Y$ , then we can take  $CZ_{21}$  to  $CY_{21}$ . Diagrammatically, this is the following argument:



The gate  $W$  that performs the required operation is

$$W = e^{i\frac{\pi}{4}X} = \begin{pmatrix} \frac{1}{\sqrt{2}} & \frac{i}{\sqrt{2}} \\ \frac{i}{\sqrt{2}} & \frac{1}{\sqrt{2}} \end{pmatrix}$$

It can be verified that  $W^\dagger Y W = Z$  and therefore the gate satisfies the above diagrammatic equation.